

1 Preseismic TEC changes for Tohoku-Oki earthquake: Comparisons 2 between simulations and observations

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9 10 **Abstract**

11
12 *Heki* [2011] reported that the Japanese GPS dense network detected a precursory
13 positive anomaly of total electron content (TEC), with $\Delta\text{TEC} \sim 3$ TECU, ~ 40 minutes
14 before the Tohoku-Oki earthquake (Mw9.0). Similar preseismic TEC anomalies were
15 also observed in the 2010 Chile earthquake (Mw 8.8), 2004 Sumatra-Andaman (Mw
16 9.2) and the 1994 Hokkaido-Toho-Oki (Mw 8.3). In this paper, we apply our improved
17 lithosphere-atmosphere-ionosphere coupling model to compute the TEC variations, and
18 compare the simulation results with the reported TEC observations. For the simulations
19 of Tohoku-Oki earthquake, we assume that the stressed associated current started ~ 40
20 minutes before the earthquake, linearly increased, and reached its maximum magnitude
21 at the time of the earthquake main shock. It is suggested that a dynamo current density
22 of ~ 25 nA m⁻² is required to produce the observed $\Delta\text{TEC} \sim 3$ TECU.

23 24 **1. Introduction**

25
26 The searching for earthquake precursors has been conducted for several decades.
27 Scientists seek for seismo-related signatures in the atmosphere or ionosphere, and
28 clarify possible signatures for precursor. Insufficiency in observational evidence drives
29 more interdisciplinary investigations attempting to unveil possible clues related to
30 earthquake activities. Several measurement methods, including VLF/LF
31 electromagnetic wave anomalies [*Hayakawa et al.*, 2010; *Hayakawa et al.*, 2012],
32 thermal anomaly [*Ouzounov and Freund*, 2004; *Ouzounov et al.*, 2006; *Pulinets and*
33 *Ouzounov*, 2011], TEC (total electron content) variations [*Liu et al.*, 2000; *Liu et al.*,
34 2001; *Liu et al.*, 2004; *Zhao et al.*, 2008] were investigated. In particular, ionospheric
35 TEC anomaly was one of the possible manifestations of seismo-ionosphere coupling
36 process [*Pulinets and Boyarchuk*, 2004; *Pulinets and Ouzounov*, 2011]. *Zhao et al.*
37 [2008] and *Liu et al.* [2009] reported that the TEC may have anomalously decreased or
38 increased up to 5 - 20% several days before the 2008 Wenchuan earthquake (Mw7.9).

39 Recently, [Heki, 2011] found that ~40 minutes before the 2011 Tohoku-Oki earthquake
40 (Mw9.0) the Japanese Global Positioning System (GPS) dense network GEONET
41 detected clear precursory positive anomaly in TEC. Similar preseismic TEC anomalies
42 were also observed in the 2010 Chile (Mw 8.8), 2004 Sumatra-Andaman (Mw 9.2) and
43 the 1994 Hokkaido-Toho-Oki (Mw 8.3) earthquakes. The finding of TEC variations
44 over the earthquake epicenter lacks physical mechanisms to explain these pre-
45 earthquake ionospheric signatures.

46

47 *Kuo et al.* [2011, 2013] proposed an electric coupling model for the lithosphere-
48 atmosphere-ionosphere (LAI) current system, as illustrated in **Figure 1**. The lithosphere
49 dynamo in the earthquake preparation region drives the internal current (\mathbf{J}_a) downward,
50 leading to the presence of a charge dipole. *Freund* [2010] has demonstrated that stressed
51 rock can generate the currents and serves as a current dynamo in the lithosphere. Due
52 to the finite conductivities in the lithosphere, atmosphere and ionosphere, the current
53 flows downward from the ionosphere, through the atmosphere (\mathbf{J}_1) and the lithosphere,
54 into the negative pole of the dynamo region. The current flowing out of the ionosphere
55 will reduce the positive charges in the ionosphere which have a higher electric potential.
56 The currents flowing in the atmosphere are obtained by directly solving the current
57 continuity equation $\nabla \cdot \bar{\mathbf{J}} = 0$ [*Kuo et al.*, 2013]. The current obtained in the
58 atmosphere can be used to calculate the electric fields at the lower boundary of the
59 ionosphere. These external electric fields are then imposed as the boundary condition
60 for the SAMI3 ionosphere model. The $\mathbf{E} \times \mathbf{B}$ plasma motion leads to TEC variations in
61 the ionosphere.

62

63 In the present study, we use this LAI coupling model to obtain the ionospheric TEC
64 variations 40 minutes before the 2011 Tohoku-Oki earthquake. Our modeling results
65 are then compared with the observed ionospheric precursor signatures (TEC variations).

66

67 **2. The Japanese GEONET TEC observation for the 2011 Tohoku-Oki Earthquake**

68

69 With the aid of the Japanese dense GPS observation network of GEONET
70 (<http://www.gsi.go.jp>), a possible anomaly for earthquake precursor could be detected
71 for the March 11, 2011, Tohoku-Oki earthquake (Mw9.0) [Heki, 2011; Heki and
72 Enomoto, 2013]. Heki [2011] used GPS-TEC data to find a clear precursory positive
73 anomaly of ionospheric TEC over the epicentral region. The TEC variations started ~40
74 minutes before the earthquake and reached nearly ten percent of the background TEC.
75 At the time of the main shock (5:46UT), eight GPS satellites were visible there [Heki,
76 2011]. The coseismic ionospheric disturbances (CIDs) can be seen by the GPS satellites

77 as the irregular TEC changes caused by acoustic waves ~10 minutes after the
78 earthquake, and the ionospheric oscillations caused by the atmospheric waves or
79 internal gravity waves 40~80 minutes after the earthquake.

80

81 **Figure 2** shows the GPS trajectories of the sub-ionospheric points (SIP) assuming a
82 thin layer at 300 km altitude. The near-by-passage of satellite 15 (red), 26 (green) and
83 27 (blue) are drawn as dots while the corresponding SIP are indicated by solid lines.
84 Here we show the detailed GPS-TEC data associated with these GPS satellites; other
85 GPS satellites have similar results. For Satellite 15, the time sequence of snapshots of
86 the geographical distribution of TEC variations are shown in **Figure 3** from UT 05:06
87 to UT 06:00 with a time step of ~5 minutes. The Japanese GEONET has more than
88 1000, and the corresponding measured Δ TEC are shown in **Figures 3, 4 and 5** where
89 each dot indicates the measured Δ TEC with color scale in units of TECU (1TECU =
90 10^{12} e/cm²) in the bottom panel.

91

92 In **Figure 3**, near the northeast side of Japan close to the west side of the 2011 Tohoku-
93 Oki earthquake epicenter, the positive anomaly of TEC is found to start at the time of
94 40 minutes before the earthquake (UT 05:46). The region with the increase of TEC
95 grew in area and reached the maximum value of Δ TEC. The TEC variations dissipated
96 and returned to normal after the CID caused by atmospheric waves generated by the
97 earthquake main shock [*Calais and Minster, 1995*].

98

99 To confirm the TEC increases preceding the Earthquake, we also show the TEC
100 measurement by Satellite 26 and 27 in **Figures 4 and 5**. The similar results of the
101 increased TEC are found; for example the covered region observed by Satellite 26 is
102 almost directly over the epicentral region. In the period from UT 05:46 to UT 05:51,
103 the observed Δ TEC can reach its peak value ~ 5 TECU. At the time of UT 06:00, it is
104 found that CID generated by earthquake main shock propagates outward, as shown in
105 the dashed circle in **Figure 4**. The oscillatory variations of the ionosphere caused by
106 atmospheric waves started at the time of ~10 minutes after the earthquake and lasted
107 40~80 minutes afterward [*Heki, 2011; Liu et al., 2011*].

108

109 **3. Simulation results from LAI coupling model**

110

111 When subjected to stress, rocks can activate positive holes (h[•]) as charge carriers and
112 generate electric currents [*Freund, 2010*]. The accumulation of positive hole charge
113 carriers at the Earth surface and charged O₂⁺ ions from field-ionization in the air near
114 the region of stressed rock. As rocks are subjected to stress, rocks activate hole (h[•])

115 charge carriers. With the exception of pure white marble, every igneous and high-grade
116 metamorphic rock tested has produced hole (h') charge carriers when stressed. The
117 positive (h') charge carriers can spread through any less stressed and even nominally
118 unstressed rock. The unstressed rock becomes positively charged while the stressed
119 rocks are negatively charged due to the loss of (h') charge carriers in the stressed region.
120 Even in oceanic region, e.g., the 2011 Tohoku-Oki earthquake in our case, charge
121 carriers have higher mobility in the ocean than in the land because of its higher
122 conductivity. The accumulated surface charge over land or ocean would drive the
123 current outward. After the charge neutralization time, some surface charges are
124 transported into the ionosphere. The equivalent effect is the current flowing into the
125 ionosphere. The direction of dynamo current flowing in the atmosphere depends on the
126 sign of the generated charges over Earth's surface near stressed rock region: downward
127 to (upward from) negative (positive) surface charge regions.

128

129 *Kuo et al. [2013]* improved the coupling model of LAI system over the previous model
130 [*Kuo et al., 2011*] which is valid only for magnetic latitude 90° and underestimates the
131 imposed electric field at the lower boundary of ionosphere. In the new model, we
132 calculate currents in the atmosphere by directly solving the current continuity equation,
133 $\nabla \cdot \mathbf{J} = 0$. The currents in the atmosphere can be solved for any arbitrary angle of
134 magnetic field, i.e., any magnetic altitude. The dynamo current density required to
135 generate the same amount of TEC variation is found to be smaller by a factor of ~ 30
136 compared to that obtained in our previous model. The typical value of dynamo current
137 J_{\max} used in the calculations is $10\text{-}100 \text{ nA m}^{-2}$, corresponding to ΔTEC of $1\text{-}7 \text{ TECU}$
138 for the daytime ionosphere.

139

140 We use the electric coupling model [*Kuo et al., 2011; Kuo et al., 2013*] to study the TEC
141 increases before the 2011 Tohoku-Oki Earthquake [*Heki, 2011*]. The simulation results
142 in our coupling models are compared with the observed TEC from GEONET. The
143 parameters in the atmosphere-ionosphere coupling model are listed below. The details
144 in the atmospheric current model and the ionosphere model are described in Section 3,
145 respectively.

146

147 **3.1. The atmospheric current model**

148

149 Our assumed atmospheric current model:

150

- 151 ■ Fault region: 450 km in length and 200 km in width [*Heki, 2011*], azimuth angle
152 ~ 30 degree from North

- 153 ■ Shift 1.5° west in longitude for EQ epicenter (38.3N,142.4E) toward the land
 154 ■ Maximum current density $J_{\max} = 25 \text{ nA m}^{-2}$
 155 ■ Current density linearly increasing from zero to its maximum value in the 40
 156 minute period (UT 05:06-05:46) before the main shock

157

158 In our atmospheric current model, we assume current distribution near the ground
 159 surface, which is confined to a region with the length $2a$ and the width $2b$.

160

$$161 \quad J_{\text{surf}}(x, y) = \frac{-J_{\max}}{4} \left[1 + \cos \frac{\pi(x-x_0)}{a} \right] \left[1 + \cos \frac{\pi(y-y_0)}{b} \right] \quad (1)$$

162

163 for $x_0 - a < x < x_0 + a$ and $y_0 - b < y < y_0 + b$, where the center (x_0, y_0) of charge

164 region is located near the epicenter. The negative sign in above equation indicates the
 165 current flowing downward. The maximum current density J_{\max} is 25 nA m^{-2} , and the
 166 total current can be integrated as $I = a \times b \times J_{\max}$. We assume a generated current source
 167 region with $a = 200 \text{ km}$ and $b = 450 \text{ km}$, which is about the size of the fault region for
 168 the Tohoku-Oki earthquake.

169

170 The current system in the atmosphere is numerically solved using $\nabla \cdot \mathbf{J} = 0$ in 3D
 171 Cartesian coordinates (x, y, z) where the x -axis is east-west, the y -axis is north-south, -
 172 $1000 \leq x, y \leq 1000 \text{ km}$, and the z -axis is the altitude, $0 \leq z \leq 200 \text{ km}$. The upper

173 ionospheric boundary condition is $\frac{\partial J_z}{\partial z} = 0$. **Figure 6** shows an example of dynamo

174 current with $J_{\max} = 25 \text{ nA m}^{-2}$, $a = 200 \text{ km}$ and $b = 450 \text{ km}$: **Figure 6a** for the current
 175 density in the $y = 0$ plane, and **6b** for that in the $x = 0$ plane, and the white lines indicate
 176 the current flows. The peak current density at altitude $z = 85 \text{ km}$ is about -12.5 nA m^{-2} .

177

178 The nearly upward or downward current \mathbf{J} flowing at 85 km altitude generally makes
 179 an angle with the inclined magnetic field. The imposed electric field on the lower
 180 boundary of the ionosphere can be derived by $E = \bar{\sigma}^{-1} J$ where conductivity tensor $\bar{\sigma}$
 181 is expressed by [Park and Dejnakintra, 1973],

182

$$183 \quad \bar{\sigma} = \begin{pmatrix} \sigma_1 & \sigma_2 \sin \theta_b & \sigma_2 \cos \theta_b \\ -\sigma_2 \sin \theta_b & \sigma_1 \sin^2 \theta_b + \sigma_0 \cos^2 \theta_b & (\sigma_1 - \sigma_0) \sin \theta_b \cos \theta_b \\ -\sigma_2 \cos \theta_b & (\sigma_1 - \sigma_0) \sin \theta_b \cos \theta_b & \sigma_1 \cos^2 \theta_b + \sigma_0 \sin^2 \theta_b \end{pmatrix}, \quad (2)$$

184

185 where σ_0 , σ_1 and σ_2 are the conductivity along the magnetic field, Pedersen
186 conductivity and Hall conductivity, respectively; θ_b is the inclined angle of the
187 magnetic field line and the horizontal plane. The values of the elements of $\vec{\sigma}$ are
188 adopted from ionosphere model SAMI3 (see below). **Figure 6c** shows the imposed
189 electric field on the upper (lower) boundary of the atmosphere (ionosphere) for the
190 current distribution in **Figures 6a and 6b**.

191

192 The imposed electric field at the lower boundary of ionosphere can be used to study the
193 TEC variations. The conductivity along the magnetic field-of-line in the ionosphere is
194 very high. The potential along the field-of-line is nearly equal potential. The imposed
195 electric field can change the electric field potential along the field-of-line in the
196 ionosphere. Therefore, we impose the electric field caused by the upward current from
197 the lower atmosphere, which is served as the electric disturbance source in the
198 ionosphere.

199

200 **3.2. The ionosphere model coupling with atmospheric current system**

201

202 The parameters in the ionosphere model (SAMI3) are:

203

- 204 ■ Day 70 (Mar 11) in 2011
- 205 ■ Solar photoionization in the ionosphere (TEC)
206 F10.7 index =150, and F10.7A=150 (81-day average of the daily F10.7)
- 207 ■ Geomagnetic Disturbance Index
208 AP =4 (mild geomagnetic condition)
- 209 ■ Neutral wind model: HWM07
- 210 ■ Simulation region +/- 8° in longitude, grid size (nf, nz, nl)=(240,101,70)

211

212 The NRL three-dimensional ionosphere simulation code SAMI3
213 (<http://wwwppd.nrl.navy.mil/sami2-OSP/index.html>), including ion dynamics and
214 electric potential, is used to investigate the TEC variation caused by the electric field
215 from the source charge of earthquake fault zone. We solve the current continuity
216 equation ($\nabla \cdot \mathbf{J} = 0$) in the ionosphere [Huba *et al.*, 2008; Huba *et al.*, 2009a; Huba *et*
217 *al.*, 2009b; Huba *et al.*, 2009c], and obtain the electric potential in the ionosphere model
218 SAMI3. The resulting electric field is used to study the plasma motion in the ionosphere
219 caused by the source charge of the earthquake fault zone.

220

221 4. Comparisons between modeling results and the observation

222

223 The 2011 Tohoku-Oki earthquake had a fault region of ~450 km in length and ~200 km
224 in width along the Japan Trench where the Pacific Plate subducts beneath NE Japan,
225 as modeled above [Heki, 2011]. The orientation of the fault region has an azimuth angle
226 ~30 degree from north centered at epicenter (38.3N, 142.4E). It is assumed that the
227 maximum current density $J_{\max} = 25 \text{ nA m}^{-2}$ increases linearly from zero to its maximum
228 value in the 40 minute period (UT 05:06-05:46) before the main shock, as shown in
229 **Figure 7**, since the increase of TEC is found to start at the time of 40 minutes before
230 the earthquake (UT 05:46), and the region with the increase of TEC grew in area and
231 reached the maximum of ΔTEC .

232

233 4.1. Simulation results of currents from the atmosphere

234

235 In comparison with the simulation results of *Kuo et al.* [2013], the modeling results
236 show the presence of the eastward (westward) electric field for downward (upward)
237 dynamo current flowing from the atmosphere into the ionosphere. At magnetic latitude
238 30° , close to the epicenter, the imposed eastward (westward) electric field causes the
239 nearly upward-northward or downward-southward direction of $\mathbf{E} \times \mathbf{B}$ motion for
240 ionospheric plasma, shown in **Figure 8a**. For the nearly upward-northward direction of
241 plasma motion with eastward electric field caused by the downward current, the $\mathbf{E} \times \mathbf{B}$
242 motion drives the ionospheric plasma from the higher density region to the lower
243 density region, enhancing the plasma density (**Figure 8c**) and increasing the TEC
244 (**Figure 8b**). Hence, we choose the downward current with eastward electric field as
245 our dynamo current.

246

247 The typical value of dynamo current J_{\max} used in the calculations is $10\text{-}100 \text{ nA m}^{-2}$,
248 corresponding to ΔTEC up to 1-7 TECU in the daytime case, shown in **Figure 9**. It is
249 also found that, in the nighttime case, the smaller value of dynamo current ($1\text{-}10 \text{ nA m}^{-2}$)
250 can lead to similar ΔTEC values. In our calculation, the dynamo current equals to the
251 multiplication of ionospheric conductivity and caused electric field. The typical
252 daytime ionospheric conductivity is ten times of the nighttime conductivity. Therefore,
253 the greater current density are required to reach the equivalent ΔTEC for the daytime
254 ionosphere.

255

256 4.2. Observation results in comparison with simulation results

257

258 **Figures 10a -10c** show the observed TEC variations at SIP of more than 1000 ground
259 GPS sites in the Japanese GEONET and their corresponding Δ TEC measurements are
260 indicated by the color dots in units of TECU. **Figures 10d-10f** show the TEC contour
261 lines from the simulation. **Figures 10e-10i** show the filled color contours of TEC where
262 the color code indicates the value of TECU. The applied eastward electric field leads to
263 the upward $\mathbf{E} \times \mathbf{B}$ motion and the increase of TEC. The Δ TEC shown in **Figures 10g,**
264 **10h and 10i** can be used to compare with measured Δ TEC results in **Figures 10a, 10b**
265 **and 10c.**

266

267 **Figure 11** shows the comparison of Δ TEC profiles from modeling results (red dots)
268 with observation (blue dots) in units of TECU one minute before the time of main shock.
269 **Figure 11a** is for the profile at geolontitude 139° , 11b at 140° and 11c at 141° . **Figures**
270 **11d, 11e and 11f** are for the profiles at geolatitude 36° , 38° and 40° . The modeling
271 results with $J_{\max} \approx 25 \text{ nA m}^{-2}$ are approximately matched with observations results.

272

273 **5. Summary and discussions**

274

275 *Heki* [2011] reported that ~ 40 minutes before the 2011 Tohoku-Oki earthquake the
276 Japanese GPS dense network detected clear earthquake precursor signals of positive
277 TEC variations over the epicentral region. We use the LAI coupling model to reproduce
278 the observed Δ TEC 40 minutes before 2011 earthquake. We assume the area of dynamo
279 current is similar to the earthquake fault region with a length $2a$ and a width $2b$ where
280 $a = 200 \text{ km}$ and $b = 450 \text{ km}$. It is found that the required dynamo current with the
281 magnitude of $10\text{-}100 \text{ nA m}^{-2}$ can produce Δ TEC of $1\text{-}7 \text{ TECU}$. In order to explain the
282 observed Δ TEC $\sim 3 \text{ TECU}$ by *Heki* [2011; 2013], the dynamo current with $J_{\max} = 25$
283 nA m^{-2} is required.

284

285 There are several areas for improvement and for future study. First, in our study, we
286 have to assume a dynamo current source. More work on the dynamo source in the
287 Earth's lithosphere is needed. The assumed dynamo current source under the ground is
288 only based on the experimental evidence of stressed rocks by *Freund* [2010] and
289 references therein.

290

291 Second, it is assumed in the SAMI3 ionosphere model that conductivity along the
292 magnetic field is infinite and the associated electric field along the magnetic field is
293 zero. In real ionosphere, we should consider the finite conductivity along the magnetic
294 field. The currents from the earthquake region flow into the ionosphere. Part of the
295 currents flow along the magnetic field, reflect from the ionosphere of the opposite

296 hemisphere, and return to the current injection region. Although our simulation results
297 show the conjugate effect, such as plasma and temperature variations, in the opposite
298 hemisphere as shown in **Figure 8**, the conjugate effect may be decreased due to a finite
299 field-aligned conductivity.

300

301 Third, it is suggested to carry out simultaneous measurements of the dynamo current
302 and electric field under the ground, the current and electric field above the Earth's
303 surface, and ionosphere TEC from ground GPS sites. The coordinated observations will
304 help to resolve the linkage among the dynamo current in the lithosphere, currents in the
305 atmosphere, and TEC variations in the ionosphere.

306

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308

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315

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317

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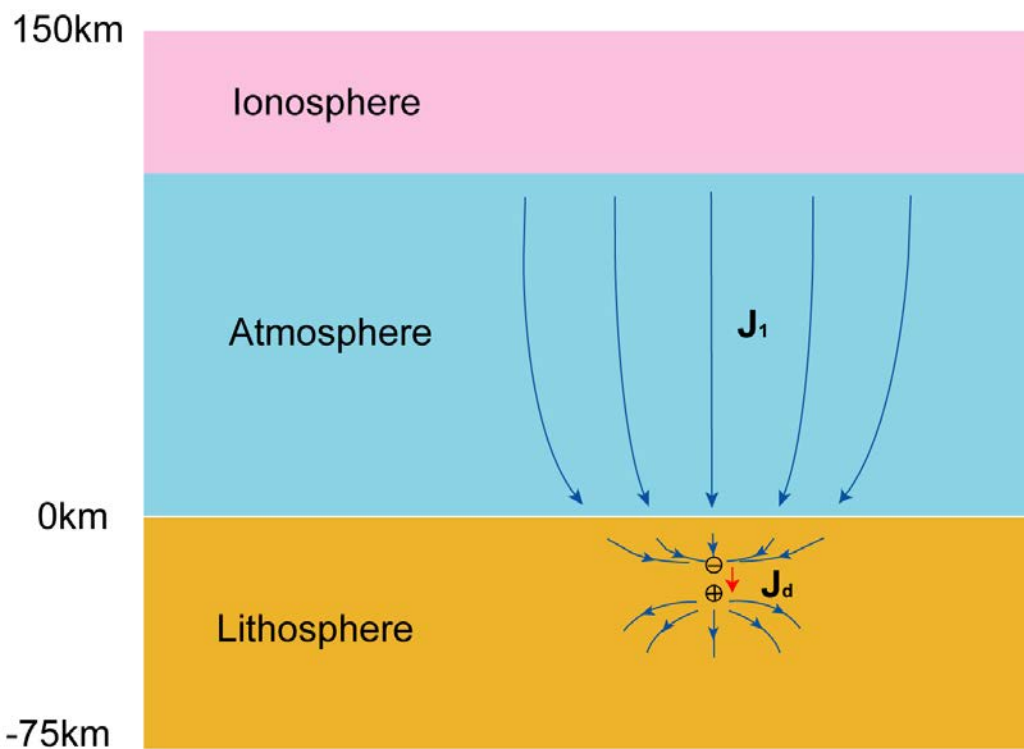
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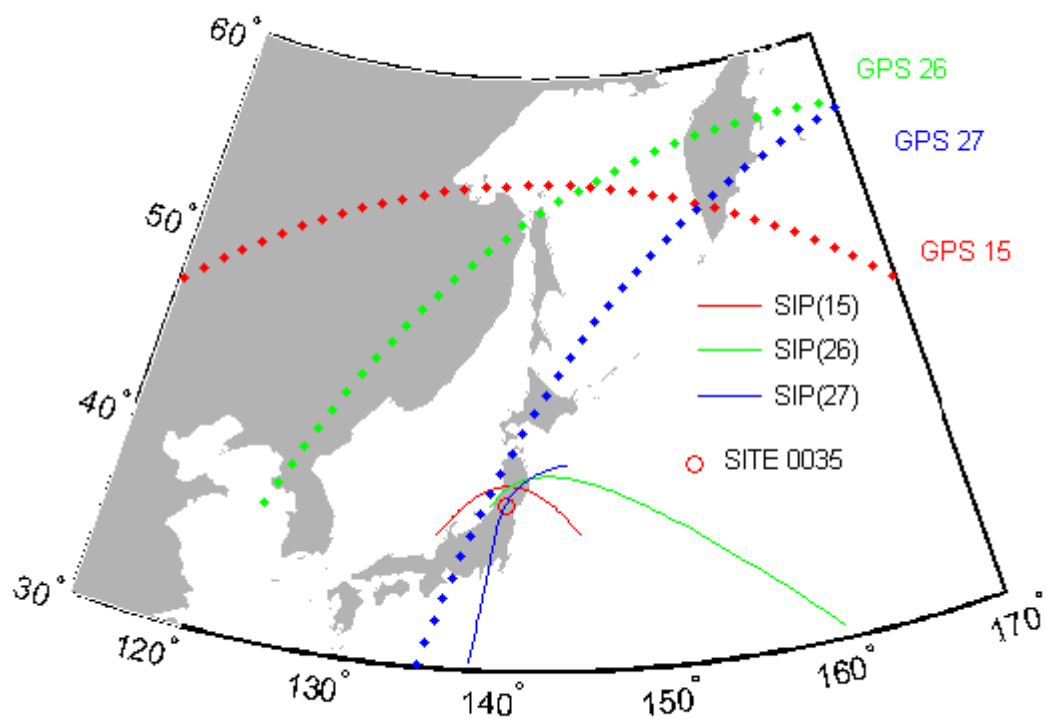
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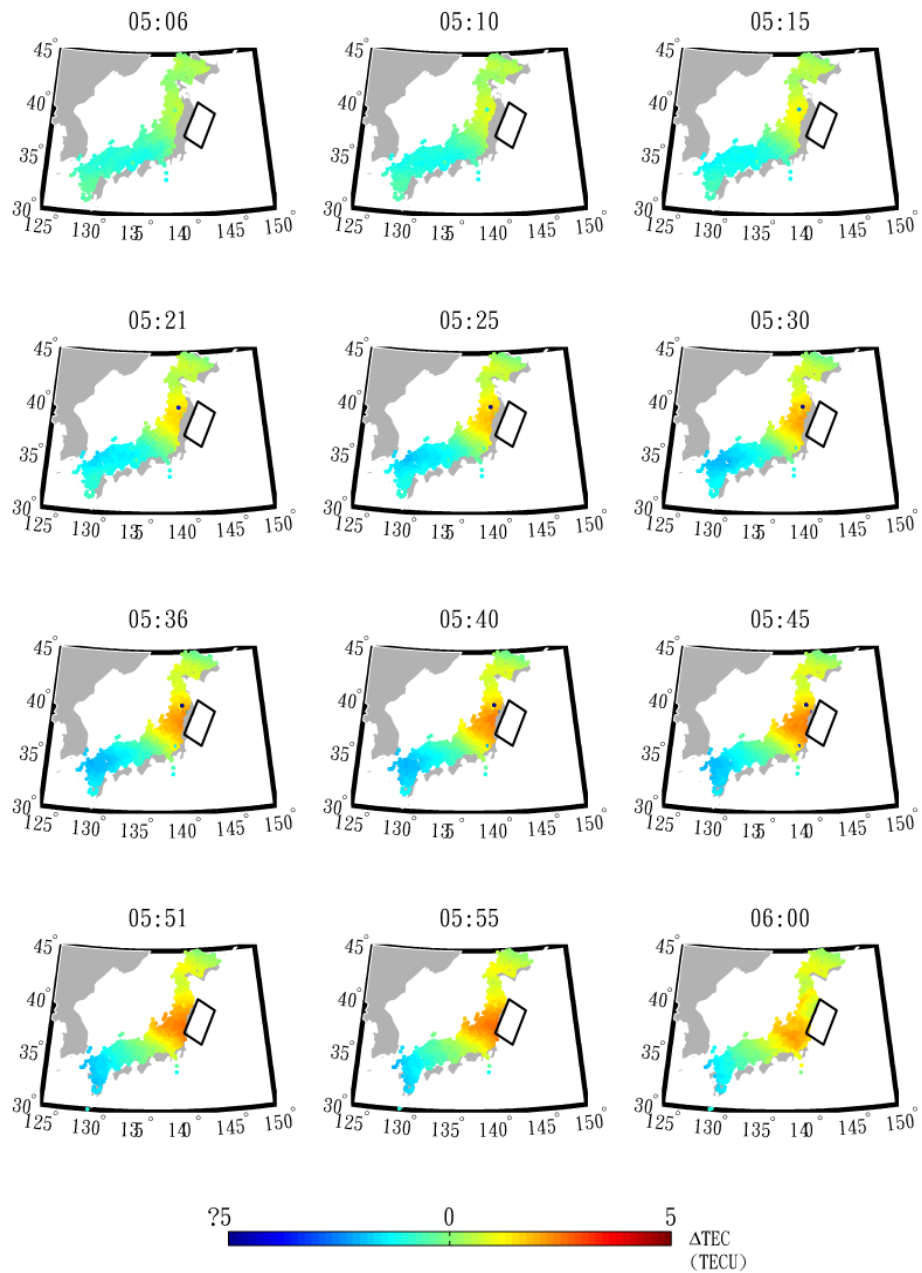
390

391 **Figure 1.** The current flow in the electric coupling model of lithosphere, atmosphere
 392 and ionosphere. The lithosphere dynamo has a charge dipole generated by the internal
 393 current J_d . The current flows downward from the ionosphere, through the atmosphere
 394 (J_1) and lithosphere, into the negative pole of the dynamo region.



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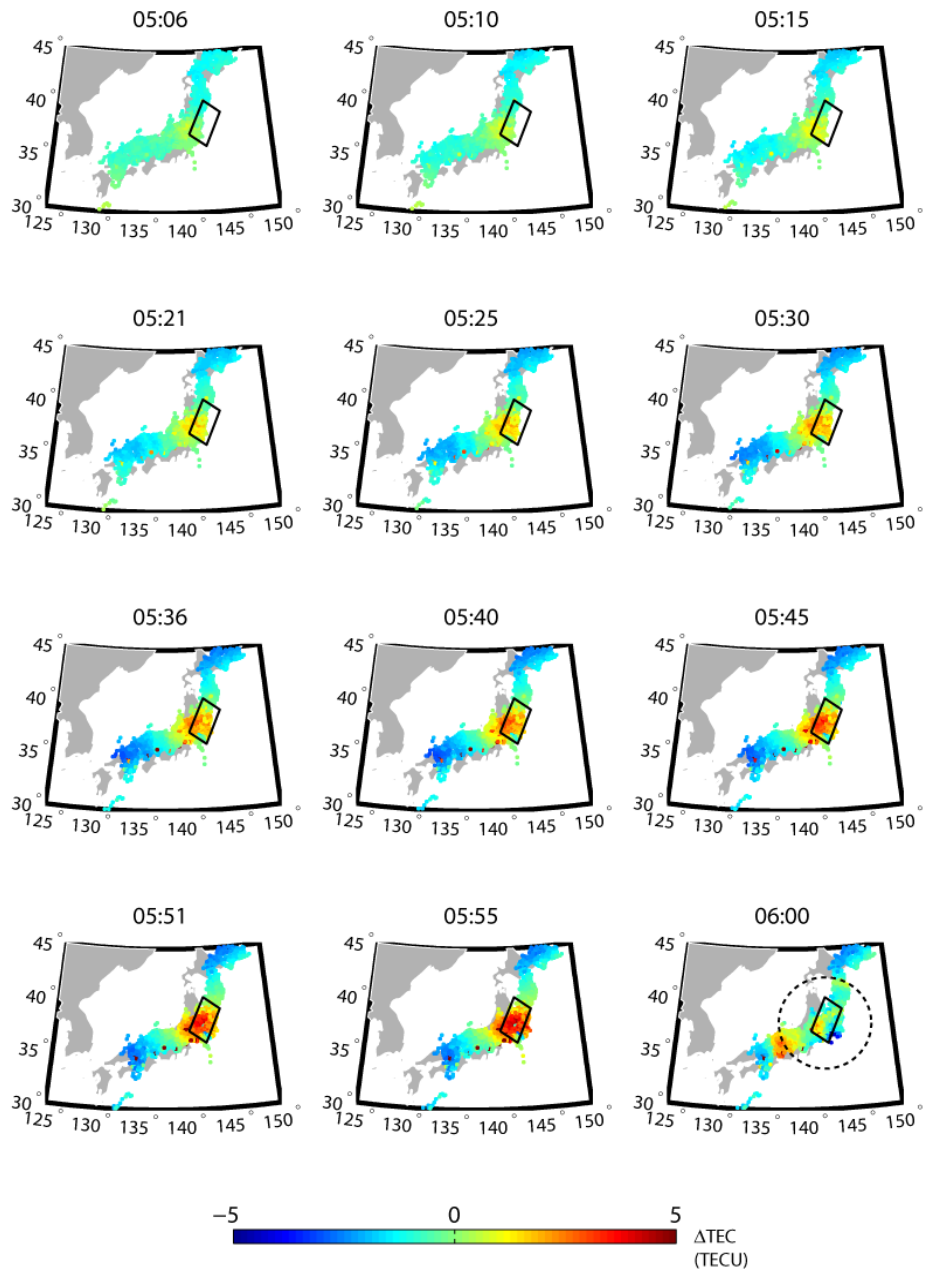
397 **Figure 2.** The trajectories of sub-ionospheric points (SIP) assuming a thin layer at 300
 398 km altitude for GPS satellites and given ground GPS site 0035. The near-by-passage of
 399 satellite 15, 26 and 27 are drawn as dots while the corresponding SIPs for GPS 15, 26
 400 and 27 are indicated by solid lines. Within the dots and solid lines, the GPS satellite 15,
 401 26 and 27 are colorized as red, green and blue lines.



402

403 **Figure 3.** The time sequence of ΔTEC recorded by the GPS satellite 15 with a time step
 404 of 5 minutes at a period 40 minutes before and 15 minutes after the 2011 Tohoku-Oki
 405 Earthquake (UT 05:46). The rectangular with black lines indicates the fault region of
 406 earthquake (~ 450 km in length and ~ 200 km in width along the Japan). The color code
 407 indicates the increase (red color) of TEC or the decrease (blue color) of TEC where the
 408 unit of TEC is TECU ($1\text{TECU} = 10^{12} \text{e}/\text{cm}^2$).

409

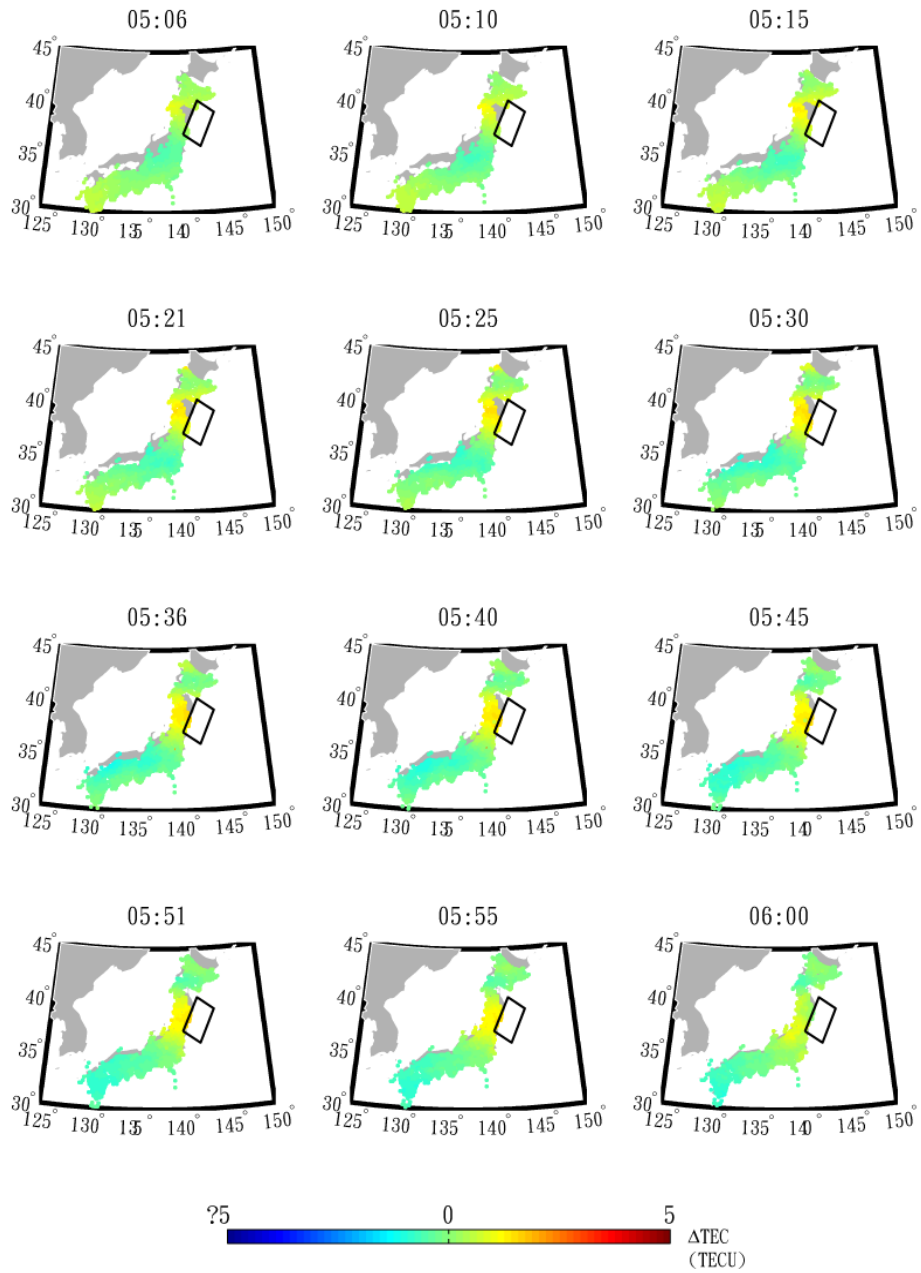


410

411 **Figure 4.** The time sequence of ΔTEC recorded by the GPS satellite 26 with a time step
 412 of 5 minutes at a period 40 minutes before and 15 minutes after the 2011 Tohoku-Oki
 413 Earthquake (UT 05:46). In the right and bottom panel, a dashed circle indicates the CID
 414 generated by earthquake propagating outwardly after the main shock.

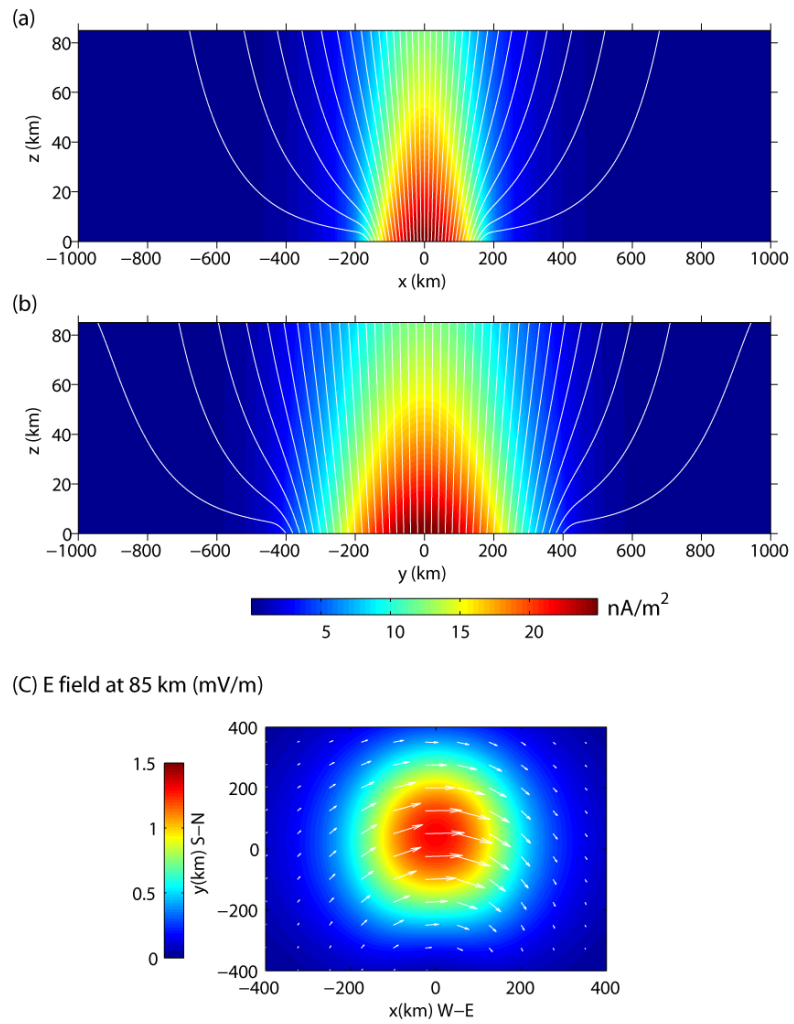
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418 **Figure 5.** The time sequence of ΔTEC recorded by the GPS satellite 27 with a time step
 419 of 5 minutes at a period 40 minutes before and 15 minutes after the 2011 Tohoku-Oki
 420 Earthquake (UT 05:46).



421

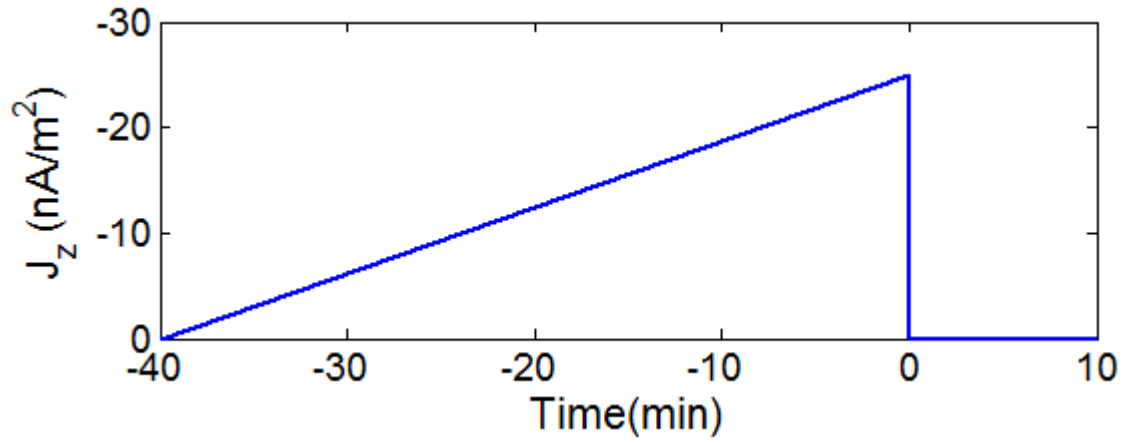
422 **Figure 6.** The distribution of current densities in (a) the $y = 0$ plane and (b) the $x = 0$

423 plane of the atmosphere. The current density is expressed in colors and the white lines

424 are current flow lines. (c) The eastward electric field at an altitude of 85 km.

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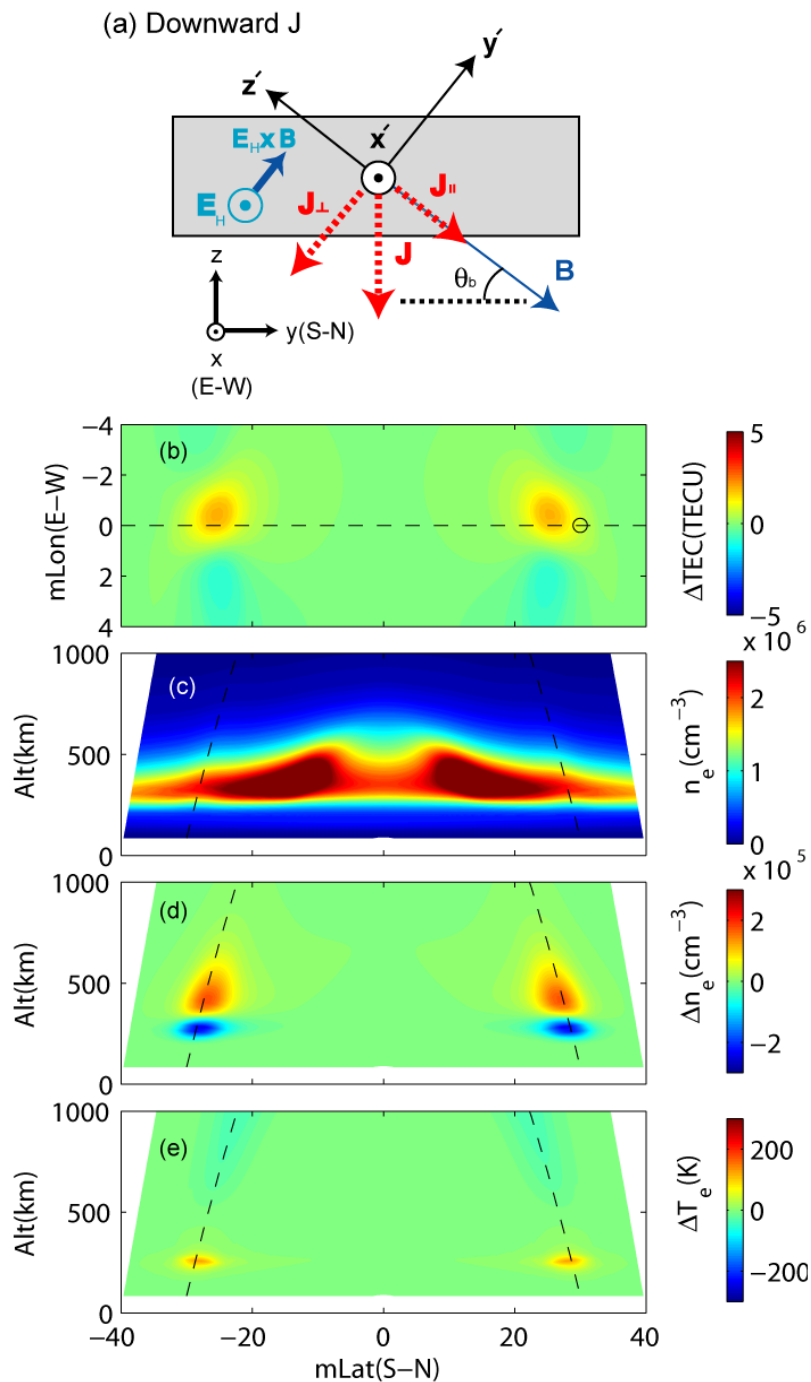
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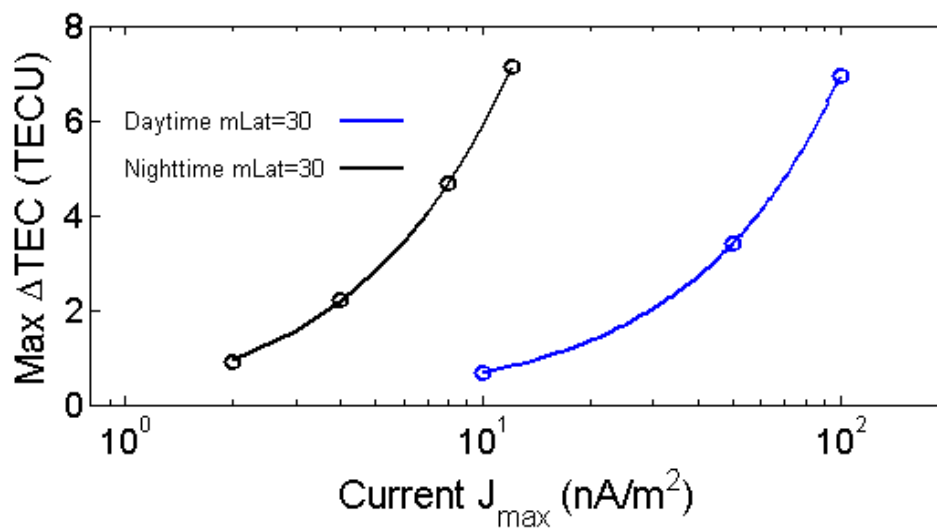
428 **Figure 7.** The maximum current density linearly increases from zero to its maximum
429 value in the 40 minute period (UT 05:06-05:46) before the main shock.

430



431

432 **Figure 8.** The ionospheric anomaly caused by downward current at the magnetic
 433 latitude 30° ; (a) the downward current lead to the presence of eastward electric field
 434 and the caused $\mathbf{E} \times \mathbf{B}$ motion enhance the ionospheric plasma density; (b) contour plots
 435 of ΔTEC in units of TECU where open circle indicate the source region; (c) contour
 436 plots of electron density n_e in the meridional planes; (d) contour plots of electron density
 437 variations Δn_e in the meridional planes; (e) temperature variations in the meridional
 438 planes.

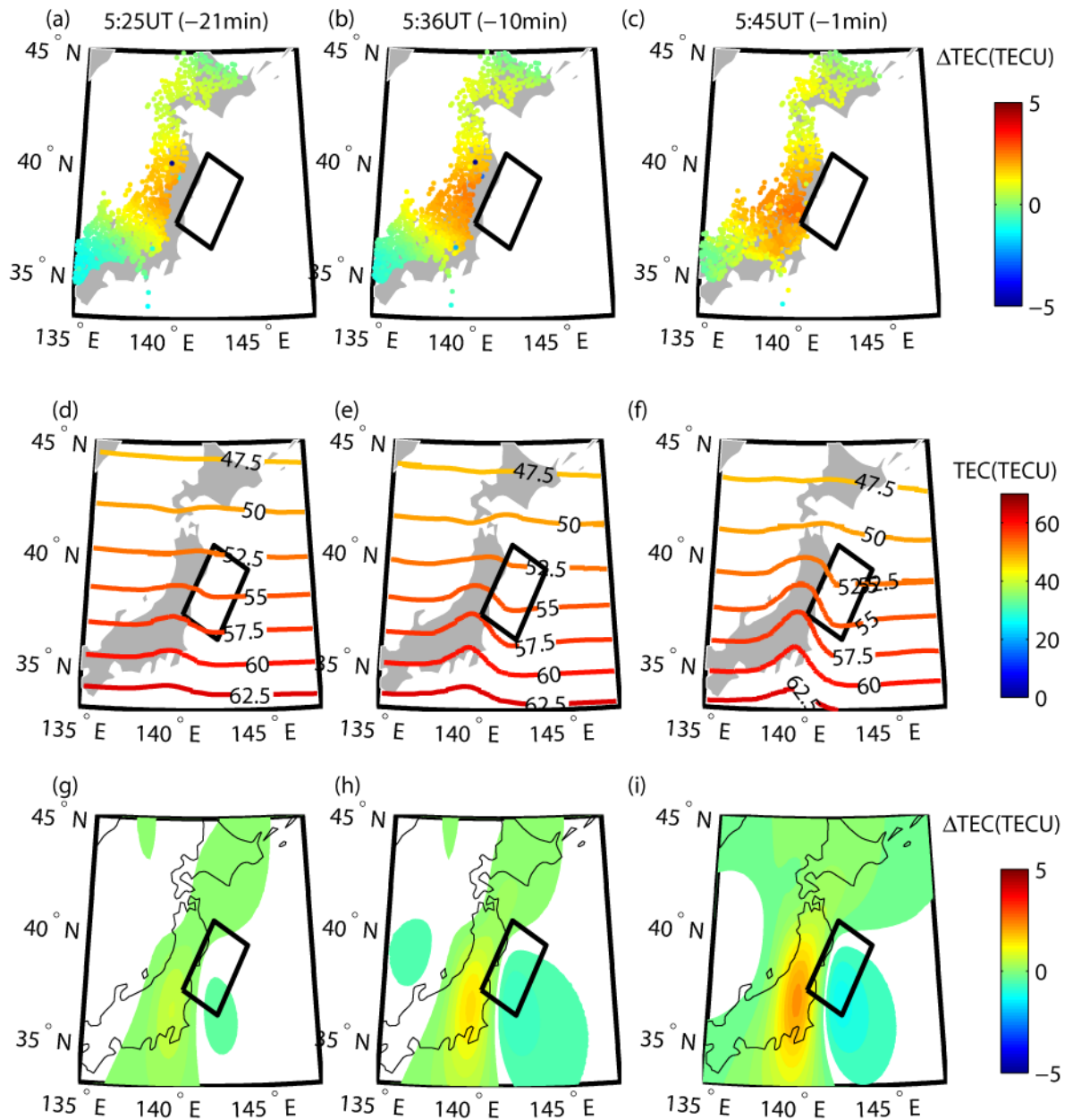


440

441 **Figure 9.** The maximum Δ TEC (TECU) varies with source current density J_{\max} in units
 442 of nA m^{-2} where the solid (dashed) lines are for Δ TEC at magnetic latitude 30° . The
 443 blue (black) lines are for daytime (nighttime) ionosphere.

444

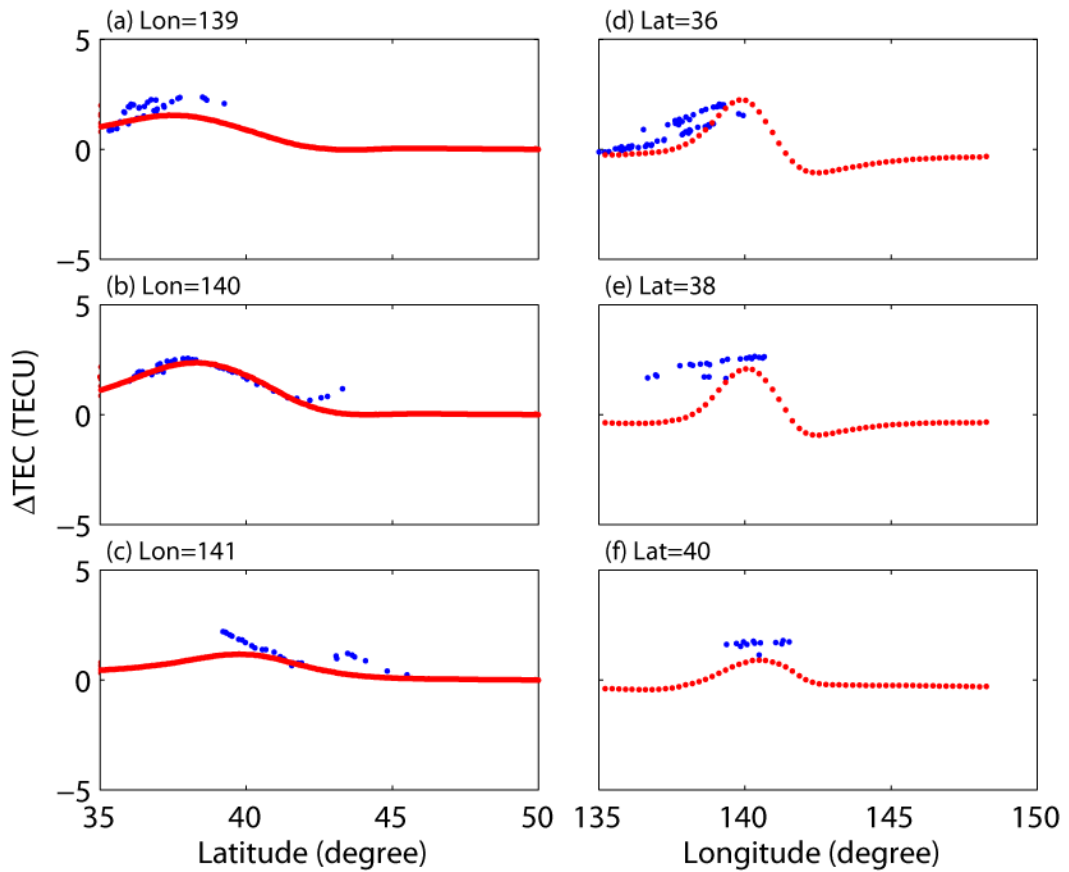
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446

447 **Figure 10.** The observed results of ΔTEC from the Japanese GEONET where color
 448 code indicates the magnitude of TEC in a time sequence of (a) 21 minutes, (b) 10
 449 minutes and (c) 1 minute before the main shock of the earthquake. The corresponding
 450 TEC contour lines from our simulation results are plotted in (d), (e) and (f). The
 451 corresponding ΔTEC from our simulation results are plotted in (g), (h) and (i).

452



453
 454 **Figure 11.** The comparison of modeling results (red dots) with observed Δ TEC (blue
 455 dots) in units of TECU at UT 05:45, one minute before the time of main shock: (a) the
 456 profile at geolontitude 139° , (b) the profile at geolontitude 140° , (c) the profile at
 457 geolontitude 141° , (d) the profile at geolatitude 36° , (e) the profile at geolatitude 38° ,
 458 and (f) the profile at geolatitude 40° .